Coaxial Cable — The Neglected Link

Is a better grade of coaxial cable worth the price difference? This analysis of the importance of shielding in coaxial lines explains why the answer is "Yes!"

By Charles Brainard,* WA1ZRS, and Ken Smith*



hen energy is transmitted through a coaxial cable, some of this energy may escape. Conversely, the cable may be in a field of energy and some of that energy can penetrate into the cable. The transformation of the energy out of, or into, the cable is called radiation, and the associated transducer is termed an antenna. The cable is a transmitting antenna when the energy escapes (egressive signals) and a receiving antenna when energy penetrates into the cable (ingressive signals). This phenomenon has many names associated with it such as leakage, radiation, isolation, shielding, shielding effectiveness, screening and screening effi-

Energy must pass through the cable shield for either egressive or ingressive signals. Obviously, for the most common use of coaxial cable, high attenuation of the energy passing through the shield is essential. There are cables that are designed for controlled leakage, and they are used when a highly controlled radiation of signals is desired. The following discussion, however, will be directed toward shields that are designed to prevent this leakage.

In 1960, a development program was established to investigate methods of construction of flexible coaxial cable to reduce cost, decrease attenuation and to

improve rf shielding. It was found that test methods used to measure the radiated energy were quite difficult and were not sufficiently repeatable to evaluate different cables. This method measured the energy external to the cable resulting from a known energy level within the cable. Hence, a new test method was developed which gave relative ratings for cables in decibels. This method was found to be useful for evaluation of different cables even though these ratings could not be directly interpreted in radiation. The repeatability of the test was in the neighborhood of 3 to 5 dB, and it was found that relative ratings varied from 20 to 100 dB for various cables.

The theoretical development of electromagnetic field coupling through the shields of coaxial cable began many years ago and the general theory was presented in an article by Schelkunoff in 1934.1 He represented the coupling by a transfer impedance and developed formulas for calculating the characteristics of solid shields. He also analyzed multiple-laver shields. Since 1934, numerous people have analyzed the coupling mechanisms and methods of measuring coupling. In the photograph, coauthor Ken Smith stands beside the Radiometer, an instrument that uses the triaxial test method2,3,4 for measurement of the electromagnetic field

coupling through the shield. He developed this instrument for the Times Wire and Cable Co. in 1978.

The purpose of this article is to show the transfer impedance and capacitive coupling impedance and, therefore, shielding effectiveness of different types of coaxial cables. The theory of electromagnetic field coupling and method of measurement will be reviewed, as will the measurement data of different types of cables.

Measurement

The Radiometer measures the absolute value of the transfer impedance and capacitive coupling impedance of the coaxial shield. An artist's sketch of the test setup is given in Fig. 1 showing that the cable is coaxially supported by a dielectric in the test chamber, creating a triaxial transmission system. The inner coaxial-transmission system is inside the test specimen and the outer coaxialtransmission system center conductor is the shield of the specimen. Its outer conductor is the test chamber. The specimen is terminated in its characteristic impedance by load A and the combination of the sweep oscillator and preamplifier. Load B and the detector are connected to the outer system by coaxial terminals. The rectangular termination on the ends of the chamber have ferrite toroids surrounding the test sample. These toroids minimize current flow along the shield of the test

^{*}Times Wire and Cable Co., 358 Hall Ave., Wallingford, CT 06492

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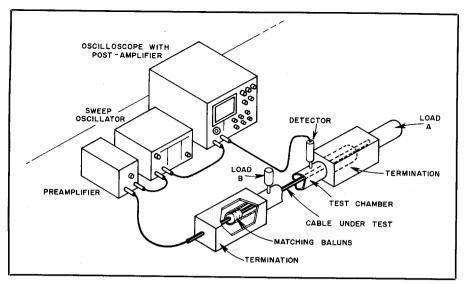


Fig. 1 — A setup for using a triaxial-test method for measuring the electromagnetic-field coupling through the shield.

specimen to the end of the rectangular termination where the shield of the specimen is grounded. These rectangular terminations form baluns creating a high impedance allowing load B and the detector to match the impedance of the chamber. Errors are not introduced by leaky conductors; the shield of the specimen is unbroken through the entire length of the fixture. The connectors on the sample are connected to the ends of the rectangular terminations and are not critical since the baluns isolate the connector leakage from the signal in the test chamber.

When the equipment is set up as shown in Fig. 1, an analysis (neglecting attenuation and assuming the cable shield is uniform) shows that the magnitude of the output voltage in the triaxial transmission system is:

$$\begin{vmatrix} V_f | = \left| \frac{(Z_t - Z_f) V_i \sin [(B_s - B_c) L/2]}{Z_s (B_s - B_c)} \right|$$
where (Eq. 1)

V_f = the detector voltage with set up of Fig. 1.

 V_i = the specimen input voltage.

 Z_t = the transfer impedance in ohms per meter.

 Z_f = the capacitive-coupling impedance in ohms per meter.

L = the distance between the coaxial terminals of the test chamber in meters.

 Z_s = the specimen characteristic impedance in ohms.

B_s = the specimen phase constant inradians per meter.

 B_c = the test-chamber phase constant in radians per meter.

An analysis of the output voltage with load B and detector exchanged shows:

$$|VR| = \left| \frac{(Z_t + Z_f) V_i Sin [(B_s + B_c) L/2]}{Z_s (B_s + B_c)} \right|$$
(Eq. 2)

where

VR = The detector voltage with load B and detector swapped.

 Z_t and Z_f may be calculated since ratios V_f/V_i , VR/V_i , B_s , B_c , L and Z_s can be measured. The test procedure provided with the Radiometer includes tables which can be used to convert the voltage ratios measured in dB to transfer impedance and capacitive-coupling impedance. The minimal specimen attenuation is neglected but the chamber attenuation is accounted for. The tables were obtained from the following equations:

$$Z_{t} = \frac{I}{L} \sqrt{Z_{s}} \frac{Z_{c}}{Z_{c}} \left(\left| \frac{\phi}{\sin \phi} \right| e^{x} + \left| \frac{\theta}{\sin \theta} \right| e^{y} \right)$$

$$(Eq. 3)$$

$$Z_{f} = \frac{I}{L} \sqrt{Z_{s}} \frac{Z_{c}}{Z_{c}} \left(\left| \frac{\phi}{\sin \phi} \right| e^{x} - \left| \frac{\theta}{\sin \theta} \right| e^{y} \right)$$

$$(Eq. 4)$$

where

 Z_t = the transfer impedance in ohms per meter.

Z_f = capacitive coupling impedance in ohms per meter.

Z_s = the specimen characteristic impedance in ohms.

L = chamber length in meters.

 $\phi = (B_s + B_c) L/2$ in radians.

 $\theta = (B_s - B_c) L/2$ in radians.

Z_c = the chamber characteristic impedance in ohms.

B_s = specimen phase constant in radians per meter.

B_c = chamber phase constant in radians per meter.

$$x = \frac{DBR - \alpha_{c/2}}{8.68} = Ln \ VR/Vi$$

$$y = \frac{DBF - \alpha_{c/2}}{8.68} = Ln \ VF/Vi$$

$$\alpha_c = \text{chamber attenuation}$$
(Eq. 6)

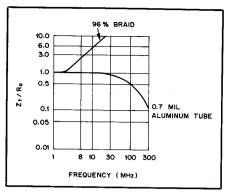


Fig. 2 — Coaxial-cable transfer impedances (Z_1) divided by the dc resistance of the shield (R_0) vs. frequency.

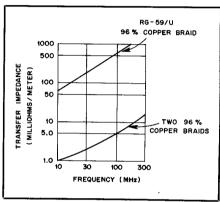


Fig. 3 — The relationship of transfer impedance vs. frequency for braided shields is shown by this graph.

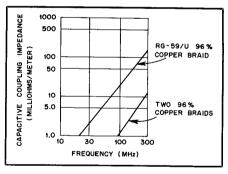


Fig. 4 — A graph representing capacitive coupling impedance vs. frequency for braided shields.

Note: X, Y, DBR, DBF and α_c are negative quantities.

Discussion and Analysis of Transfer-Impedance Test Results

Transfer impedance in an elementary length of coaxial cable is defined as the ratio of the potential gradient (voltage) in the disturbed circuit to the current flowing in the interfering circuit. When the cable is acting as a transmitting antenna (egressive signals) the disturbed circuit is the environment around the cable. When the cable is acting as a receiving antenna (ingressive signals) the disturbed circuit is

within the cable and the interfering circuit is the environment around the cable. A lower transfer impedance reduces the electromagnetic coupling (radiation).

The transfer impedance of a braided shield has two components — a diffusion component caused by current diffusing through the metal and a mutual-coupling component caused by penetration of the magnetic field through the openings in the braid. The mutual-coupling component can be represented by a mutual inductance. (Figs. 2, 3 and 4 are related to this discussion.)

The transfer impedance is the vector sum of these two complex quantities, and its magnitude is:

$$|Z_t|$$
= $\sqrt{(|Z_d|\cos\phi)^2 + (|Z_d|\sin\phi + |Z_m|)^2}$
(Eq. 7)

where

 $\phi = 0.785 - \tan^{-1} (\cot d/\delta \tan d/\delta)$ Zd = the diffusion component of Z_t in ohms per meter.

 $Zm = the mutual-coupling component of <math>Z_t$ in ohms per meter.

d = the diameter of braid wire in meters.

 δ = the skin depth in meters.

The approximate diffusion component and mutual-coupling component for braided cable is obtained from an extension of Vance's equation and Schelkunoff's.'

The diffusion component is:

$$|Z_d| = R_{dc} \frac{(\sqrt{2}) d/\delta}{\sqrt{\sinh^2 (d/\delta) + \sin^2 (d/\delta)}}$$
(Eq. 8)

$$\delta = \sqrt{\frac{\rho}{\pi f u'}}$$
 (Eq. 9)

The mutual-coupling is:

$$|Z_{\rm m}| = \frac{\omega v \, \mu \rm m}{\pi^2 \, D^2} \tag{Eq. 10}$$

where

 ρ = the resistivity of the shield in ohms per meter.

f = the frequency in hertz.

 μ' = the absolute magnetic permeability of the shield in henrys per meter.

d = the diameter of braid wire in meters.

 R_{dc} = the dc resistance of the shield in ohms per meter.

 ω = the angular frequency in radians per second = $2\pi f$.

v = the number of holes per meter in the braided shield.

μ = the absolute magnetic permeability of the insulation between the conductors in henrys per meter.

D = the mean inside diameter of the shield in meters.

m = the magnetic polarizability of the holes in the braid.5

Z_m = the mutual-coupling component in ohms per meter.

Z_d = diffusion component in ohms per meter.

 $\delta = \text{skin depth in meters.}$

Analysis of Capacitive-Coupling Impedance Test Results

The openings in the shield also allow the electric field to penetrate, creating electric coupling. This coupling can be represented by a capacitive coupling between the center conductor of the coaxial cable and the return path external to the cable.

The capacitive-coupling impedance is derived for the definition accepted by the International Electrotechnical Commis-

sion Working Group I (Screening Efficiency) and Vance's equation for transfer admittance.

$$Z_{\rm f} = \frac{p}{m} Z_{\rm m} \sqrt{\epsilon r_{\rm c}/\epsilon r_{\rm i}}$$
 (Eq. 11)

where

Z_f = the capacitive coupling impedance in ohms per meter.

p = the electric polarizability of the holes in the braid.

m = the magnetic polarizability of the holes in the braid.

 Z_m = the mutual-coupling component of Z_t in ohms per meter.

 ϵr_e = the relative dielectric constant of the insulation in the external circuit.

 ϵr_i = the relative dielectric constant of the insulation within the cable.

The capacitive-coupling impedance will be zero if there are no openings in the shield. If openings exist, then the capacitive-coupling impedance should vary directly with frequency. The test data plotted in Fig. 4 follows this characteristic reasonably well.

Conclusions

Measurements of the transfer impedance and capacitive-coupling impedance of coaxial shields can be made. Results agree with theoretical equations. Since the theory of transfer of energy through shields is known, an engineer can analyze and design coaxial cable theoretically. Because of the different types of coaxial cables in use today, the design engineer should be aware of the large variation in the coupling of electromagnetic fields through the shields.

The information in Table 1 gives an idea of the relative isolation of coaxial cable in accordance with percent and type of shielding. We have also indicated the losses in dB per 100 feet at 15 and 150 MHz (near two very popular Amateur-Radio bands) for each construction in the table. While Table 1 deals with RG-59/U cable, the isolation characteristics are applicable to any solid-dielectric cable while the losses in dB per 100 feet can be interpolated to other coaxial lines.

Down-line attenuation of cables is controlled by (1) the type and material of the center conductor; (2) the velocity of propagation and type of material of the center dielectric; and (3) the type and material of the outer conductor. An examination of the loss curves of Fig. 5 shows that RG-58 A/U exhibits loss characteristics from 11% to 12% more than that of RG-58/U. The only difference between these two cables is the nature of the center conductor. RG-58 A/U has a 19-strand tinned-copper center conductor. Both the stranding and tinning increase attenuation. Tinned copper is very popular with antenna manufacturers because it simplifies soldering. But, considering attenuation tinned-copper shields

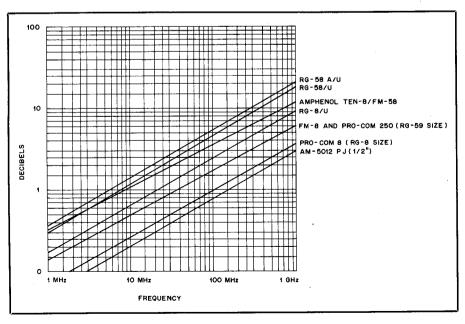


Fig. 5 - Nominal loss characteristics per 100 ft (30.48 m).

Table 1 Relative isolation characteristics of RG-59/U coaxial cable vs. percentage and type of braid coverage/shielding.

Shield	Relative Isolation (dB)	Ratio of Power Radiated from Cable	Losses in dB per 100 ft. (30.48 m)	
			15 MHz	150 MHz
40% bare copper	17	1:50	1.72	5.55
51% bare copper	18	1:63	1.72	5.55
59% bare copper	26	1:398	1.39	4.51
79% bare copper	34	1:2500	1.13	3.67
98% bare copper	52	1:160,000	0.98	3.20
96%/96% bare copper	83	1.2×10^{8}	1.01	3.31
Solid sheath (alum.)	282	$1:17 \times 10^{28}$	0.89	2.91

Note: Isolation capabilities of coaxial cable at 20 meters is roughly 10 times as good as at 2 meters.

Table 2 Formulas Common to All Coaxial Cable

Capacitance (C) =
$$\frac{7.36E}{\text{Log (D/d)}}$$
 picofarads/ft

Inductance (L) = 0.140 Log (D/d) microhenrys/ft

Impedance (Z_o) =
$$10^3 \sqrt{\frac{L}{C}} = \frac{138}{\sqrt{E}} \text{ Log}$$

(D/d) ohms

Velocity of propagation as % of speed of light =
$$\frac{100}{\sqrt{E}}$$

Time delay = $1.016 \sqrt{E}$ nanoseconds/ft

Cutoff frequency =
$$\frac{7.50}{\sqrt{E(D+d)}}$$
 = F_{co} (GHz)

Magnitude of Reflection Coefficient =

$$(1) = \left[\frac{Z_{r} - Z_{o}}{Z_{r} + Z_{o}}\right] = \frac{VSWR - 1}{VSWR + 1}$$

$$VSWR = \frac{1 + \Gamma}{1 - \Gamma}$$

*Varies with frequency

Peak voltage = $\frac{1.15 \text{Sxd} (\log D/d)}{K}$

$$\alpha = \frac{0.435}{Z_0(D)} \left[\frac{D}{d} K_1 + K_2 \right] \sqrt{F} + 2.78$$

$$\sqrt{E}(P,F,) (F)$$

where

- α = attenuation in db/100 ft
- d = outside diameter of inner conductor
- in inches
- D = inside diameter of outer conductor in inches
- = maximum voltage gradient of the cable insulation in volts per mil
- = relative dielectric constant of the insulation of cable

Log = Logarithm to base 10

K = safety factor

K₁ = strand factor and material

K2 = braid factor and material

= frequency in MHz

P.F. = power factor

Feet x 0.3048 = meters

generate seriously detrimental effects above 500 MHz.

Another example that can be drawn from Fig. 5 is loss differentials (approximately 33%) between RG-8/U and FM-8 (Flexifoam by Times Wire), RG-8/U utilizes a 7-strand bare-copper center conductor and a 66% velocity core dielectric as opposed to a solid bare-copper center conductor and 79% velocity center core. It should be noted that in order for the FM-8 to possess a 50-ohm characteristic, the center conductor must be enlarged to 0.102 inches (2.59 mm) or no. 10 gauge vs. 0.0808 inches (2.05 mm) or no. 12 gauge for the center conductor on RG-8/U. As the velocity of propagation of a coaxial cable is increased while maintaining the same inside diameter of the shield, the center conductor must be increased in size in order to maintain the same cable impedance.

Braided-shield coaxial cables are, to varying degrees, impacted by their environment. That is, mounting a poorly shielded cable directly to a tower leg, as is so common, can drastically alter the attenuation characteristics of the cable. This change in characteristics is often in excess of 10 times. Degradation increases with frequency.

Placing poor cables in any conducting environment such as when attached to a tower leg or even buried in the ground can cause adverse results.

From the formulas in this article, transfer impedance and capacitivecoupling impedances can be calculated. As these impedance values rise, the outer conductor (shield) has larger openings. Consequently, higher-impedance cable is more affected by environment.

There is much more to coaxial cable than meets the eye. Since the cost of coaxial transmission line is usually relatively small in proportion to other station costs, it's difficult to reason the use of poor grade cable. A properly shielded line should be a must for all installations.

Properties of Wire and Cable Insulating Materials

Material	Dielectric Constant	Power Factor	Volume Resistivity (ohms-cm)	Operating Temperature Limits (°C)
TFE Polyethylene Cellular polyethylene Polyvinylchloride Nylon Kel-F Silicone rubber Ethylene propylene FEP Perforated TFE Cellular FEP Polyimide	2.1 2.3 1.40-2.10 3.00-8.00 4.60-3.50 2.37 2.08-3.50 2.24 2.10 1.50 1.40 1.50 3.00-3.50	0.0003 0.0003 0.0003 0.0700-0.1600 0.040-0.030 0.0270-0.0053 0.007-0.016 0.00046 0.0003 0.0002 0.0002 0.0002 0.0002*	1019 1016 1012 2 × 1012 4 × 1014 1.2 × 1018 1013 1017 1018 1019 1019 1818 1013	-75 + 250 -75 + 80 -75 + 80 -55 + 105 -60 + 120 -40 + 150 -70 + 250 -40 + 105 -75 + 250 -75 + 250 -75 + 250 -75 + 300
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Normal

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